# Hybrid NMR: A Union of Solution and Solid-State NMR 

Supporting Information

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Distribution of crystal orientations and symmetric tensors. We were able to characterize the distribution in polar angles $(\beta)$ for the crystallite orientations with respect to the laboratory frame from the ${ }^{2} \mathrm{H}_{2} \mathrm{O}$ quadrupolar powder lineshape. Based on the fit of the PRO ${ }^{2} \mathrm{H}$ spectrum in Figure 1, we found that our samples are uniformly and spherically distributed in the $\beta$ angle. However, the ${ }^{2} \mathrm{H}_{2} \mathrm{O}$ quadrupolar tensor is expected to be axially symmetric, and this spectrum cannot be used to determine the crystallite distribution in the azimuthal angle, $\alpha$.

The average tensor measured in a uniformly distributed PRO sample is described as follows (from equation (16)):

$$
\begin{equation*}
V_{20}^{(L A B)}=\sum_{m} V_{2 m}^{(A V E)} \mathcal{D}_{m 0}^{A L} \tag{S1}
\end{equation*}
$$

The final spectrum is integrated over the crystallite $\beta^{\text {AL }}$ and $\alpha^{\text {AL }}$ angles from the five Wigner components, $\mathcal{D}_{m 0}^{A L}$. In the more general case, the $\beta$ and $\alpha$ angles may deviate from a uniform spherical distribution if the crystallites are locked in the hydrogel in a non-uniform distribution.

$$
\begin{equation*}
V_{20}^{(L A B)}=\sum_{m} V_{2 m}^{(A V E)} \mathscr{D}_{m 0}^{A L} P(\alpha) P(\beta) \tag{S2}
\end{equation*}
$$

This deviation is represented by a probability distribution, $P(\alpha)$ and $P(\beta)$, for the two angles. This expression assumes that the distribution functions in crystallite $\beta$ and $\alpha$ angles are not correlated.

We were able to avoid a non-uniform distribution, at least in the $\beta$ angle, by reprocessing our samples. However, it is not possible, without rotating the sample about a vector orthogonal to the $B_{0}$ field, to determine the crystallite distribution in $\alpha$ angles from an axially symmetric (or nearly axially symmetric) tensor, like the quadrupolar coupling
tensor of ${ }^{2} \mathrm{H}_{2} \mathrm{O}$. For an axially symmetric tensor, the $\mathrm{m}= \pm 2, \pm 1$ terms are equal to zero, and equation (S2) only depends on the $\beta$ angle.

$$
\begin{equation*}
V_{20}^{(L A B)}=\sum_{m} V_{20}^{(A V E)} d_{00}^{A L}(\beta) P(\beta) \tag{S3}
\end{equation*}
$$

Likewise, a non-uniform distribution in the azimuthal angle cannot make an axially symmetric tensor appear axially asymmetric. The above does not hold true if the distributions in $\alpha$ and $\beta$ angles are correlated such that $P(\alpha, \beta)$. However, in this case, the spectrum of symmetric tensors would all appear axially asymmetric, including the ${ }^{2} \mathrm{H}$ spectrum of ${ }^{2} \mathrm{H}_{2} \mathrm{O}$.


Figure S1. Simulated ${ }^{2} \mathrm{H}$ spectra for a biaxial liquid crystal ordered along the $\beta_{\mathrm{CL}}$ angle and statically disordered along the $\alpha_{C L}$ angle. Spectra were simulated using Simpson with 2000 REPULSION orientations with the $\beta_{\mathrm{CL}}$ angle aligned to the specified value. A RQC of 40.0 Hz was used and a quadrupolar asymmetry, $\eta$, of 0.0 (left panel) or 1.0 (right panel) were used. The spectra were apodized with 1.0 Hz of exponential line broadening before Fourier Transformation.


Figure S2. Pulse sequence for the SLF- ${ }^{13} \mathrm{C}-\mathrm{HSQC}$. The $\left\{{ }^{1} \mathrm{H}-{ }^{13} \mathrm{C}\right\}$ INEPT and reverse INEPT periods were incremented together in a separate dimension (F1), producing a $\sin ^{2}\left(\pi \mathrm{~J}_{\mathrm{CH}} \mathrm{t}_{1}\right)$ or $\sin ^{2}\left(\pi\left(\mathrm{~J}_{\mathrm{CH}}+\delta_{\mathrm{CH}}\right) \mathrm{t}_{1}\right)$ signal modulation. The pulse sequence includes a trim pulse (tr.) and gradient encoding (E) and decoding (D) pulses. The period $\Delta$ was equal to the initial $\mathrm{t}_{1}$ duration.


Figure S3. The $\chi^{2}$ fit for the asymmetry $(\eta)$ of the ${ }^{13} \mathrm{C}^{\prime}$ RCSA tensor from the SLF$\mathrm{HA}(\mathrm{CA}) \mathrm{CO}$ spectrum of PRO ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$-L-aspartate at pH 5.2 in ${ }^{2} \mathrm{H}_{2} \mathrm{O}$ sample. All other parameters ( $\delta_{\mathrm{iso}, \mathrm{C}^{\prime}}, \delta_{\mathrm{CH}}, \delta_{\mathrm{C}^{\prime}}, \beta, \gamma, \mathrm{lb}_{\mathrm{CH}}, \mathrm{lb}_{\mathrm{C}^{\prime}}$ ) were minimized for every $\eta$ point. See the Experimental Section for sample and experimental details.


Figure S4. Comparison of the experimental and simulated spectra of the $\left\{{ }^{1} \mathrm{H}^{\alpha}-{ }^{13} \mathrm{C}^{\alpha}\right\} /{ }^{13} \mathrm{C}$, 2D projection of a SLF-HA(CA)CO spectrum for the PRO ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$-L-aspartate at pD 5.2 in ${ }^{2} \mathrm{H}_{2} \mathrm{O}$. The two simulated spectra represent the best-fit spectrum with a ${ }^{13} \mathrm{C}{ }^{\prime}$ RCSA tensor asymmetry of $\eta_{C^{\prime}}=0.50 \pm 0.01\left(\chi^{2}=15408\right)$ as well as the spectrum with a fixed $\eta_{C^{\prime}}=0.0\left(\chi^{2}=19690\right)$.

Table S1. F-tests and quality of fits for the SLF-HACACO of ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$-L-aspartate in ${ }^{2} \mathrm{H}_{2} \mathrm{O}$.
$\left.\begin{array}{llll}\hline \begin{array}{l}\text { Free } \\ \text { Parameters }\end{array} & \begin{array}{l}\text { Number of free } \\ \text { parameters }\end{array} & \chi^{2} & \chi_{\text {red }}^{2}\end{array}\right]$ F-test P-value (\%)

